

THEORY OF NEUTRON FLUCTUATIONS WITH APPLICATIONS IN NUCLEAR SAFEGUARDS AND NOISE DIAGNOSTICS IN RESEARCH AND POWER REACTORS

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Abstract: Neutron fluctuations in both low power systems (“zero power reactor noise”) [1], [2] and power reactors (“power reactor noise”) [2], [3] carry important information about the system. The origin of the fluctuations is different in these two regimes, and consequently, they are described by different mathematical methods, as well as that they have different characteristics and application areas. At low power (critical assemblies and research reactors) the fluctuations are due to the branching process (multiple generation of several neutrons in fission); in power reactors it is the technological processes of random character, such as boiling of the coolant, vibration of mechanical structures etc., which induce the neutron noise. Common to them is that neutron fluctuations can be used to diagnose the state of the system, either by determining some of its parameters (reactivity coefficients, void content etc.), or by detecting beginning anomalies (excessive vibrations, flow anomalies etc.).

In this report the principles and selected applications of both zero power noise and the power reactor noise will be given. The current significance of research in power reactors is also indicated by some on-going collaborations, such as the large EU-supported co-ordinated research project, CORTEX, whose goal is the development and verification of core diagnostic methods [4], and the planned noise diagnostic option of the US DoE supported Versatile Test Reactor (VTR) Project [5]. A particular aspect of these projects is that the development and verification of calculational and experimental methods for power reactors is supported by dedicated experiments and measurements in zero-power and research reactors. In these latter, the perturbations need to be generated artificially; on the other hand, unlike in power reactors, determination of the parameters of the noise source as well as measurements of the induced neutron noise can be performed simultaneously under controlled circumstances.

Keywords: *neutron noise, reactor diagnostics, safeguards, research reactors, power reactors.*

1. INTRODUCTION

The theory and application of neutron noise diagnostics has a long and successful history [1] – [3]. The theory of zero power reactor noise, i.e. neutron fluctuations in low power systems or in samples of fissile material, is used to determine the subcritical reactivity of a reactor during start-up, or during the operation of an accelerator driven subcritical system (ADS). Another application is nuclear safeguards, i.e. detection, identification and quantification of fissile nuclear material, either for detecting illicit trafficking, or to verify the integrity and fissile content of spent nuclear fuel, before depositing it in the final repository.

The methods of power reactor noise diagnostics are used for a continuous, non-intrusive surveillance of power reactors. Both operational parameters in the stationary state, such as the moderator temperature coefficient in PWRs or the void content in BWRs, can be continuously determined and possible changes being detected, or the occurrence of abnormal processes, such as fuel and control rod vibrations, boiling in PWRs, flow anomalies in both PWRs and BWRs, or instability in BWRs can be detected in an early stage. This way they can be identified and

quantified at an early stage, such that preventive actions can be taken. In the forthcoming illustrations of these applications will be given.

2. ZERO POWER REACTOR NOISE

2. 1. Methodology

The methodology is based on probability balance equations, also called Kolmogorov or Chapman-Kolmogorov equations, or just master equations. Usually an energy-independent (“one-group”) description is used, with an infinite homogeneous medium, in order to arrive analytically tractable equations. The traditional theory deals with the temporal evolution of the probability distributions of discrete quantities, such as the probability $P(N,C,Z,t)$ of having N neutrons and C precursors in the system, and having Z detector counts between time 0 to t . Usually, only the few lowest order moments, such as the variance (second moment) or skweness (the third moment) of the quantities are sought.

One traditional application is to determine the subcritical reactivity of a system, driven by a stationary source, from the relative variance $\sigma_Z^2(t) / Z(t)$ of the detector counts $Z(t)$ measured during a time period t (the variance to mean or Feynman-alpha method). Accounting for only one (average) group of delayed neutrons, it is given as

$$\frac{\sigma_Z^2(t)}{Z(t)} = 1 + Y(t) = 1 + \varepsilon \left[A_1 \left(1 - \frac{1 - e^{-\alpha t}}{\alpha t} \right) + A_2 \left(1 - \frac{1 - e^{-\alpha_d t}}{\alpha_d t} \right) \right] \quad (1)$$

where ε is the detector efficiency, A_1 and A_2 constants, containing the first two moments of the number distribution of fission neutrons, and the delayed neutron fraction and decay constant. $Y(t)$ is the so-called Feynman Y -function, whose time dependence is determined by the parameters α and α_d , i.e. the prompt and delayed neutron α time constants, where

$$\alpha = \frac{\beta - \rho}{\Lambda} \quad (2)$$

with ρ being the reactivity and Λ the prompt neutron generation time. From a measurement of the function $Y(t)$, the reactivity can be extracted from the prompt neutron decay constant, which is determined by a curve fitting of (1) to the measured data. An illustration of the Feynman Y -function for different values of the reactivity ρ is shown in Fig. 1.

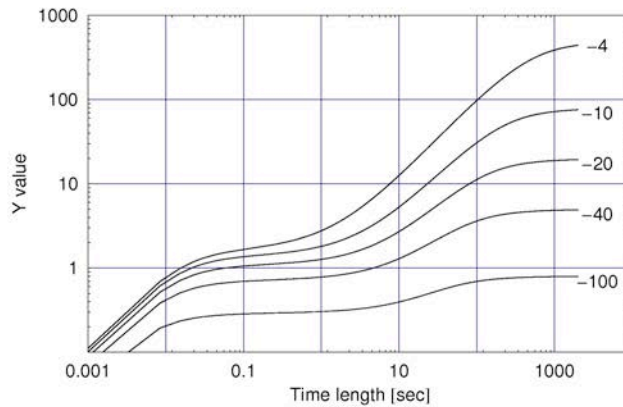


Fig. 1 Feynman-Y values as functions of time, for various values of subcritical reactivity in units of 10^{-4}

Another application is the field of nuclear safeguards. i.e. detection, identification and quantification of hidden or non-classified nuclear material. This is based on deriving theoretical formulae for the three lowest order moments of the number of detected neutrons, emitted from the item through spontaneous fission and subsequent internal multiplication, and measuring the detection rates of single, double and triple coincidences. This process makes it possible to determine a few unknowns of the sample, including sample mass. This is called “multiplicity counting”.

To give a flavor of the methodology, we show the backward master equation for the number of neutrons leaving the sample, due to one initiating neutron. The source neutron may undergo some fission reactions with probability p before escaping from the sample, hence the multiplicity (number distribution) of the neutrons per source event will differ from that of the spontaneous fission, and will be dependent on the sample mass. A backward type master equation for this probability reads as

$$p(n) = (1-p)\delta_{n,1} + p \sum_{k=1}^{\infty} p_f(k) \prod_{i=1}^k p(n_i) \quad (3)$$

$\{n_1+n_2+\dots+n_k=n\}$

and the corresponding equation for the generating function reads as

$$h(z) = (1-p)z + p q_f [h(z)] \quad (4)$$

From this generating function the factorial moments can be derived by derivation w.r.t. z .

Whereas this formalism and the applications have been known and used for quite some time, significant new developments were made presently in the methodology, which will be described below.

2. 2. New developments and results

The above methods, using discrete pulse counting, suffer from the dead time effect, and the shortage of one of the most commonly used detector materials (He-3). Hence, we have recently initiated research to see whether the same statistical information (variance of the number of counts in a reactor for reactivity measurement, or the factorial moments of neutrons detected from an unknown sample, for multiplicity measurements in nuclear safeguards) can be extracted from the statistical moments of the continuous currents of fission chambers.

This required to set up master equations for the continuous stochastic processes, represented by the detector signals, as functions of the underlying discrete detection process. Elaboration of such a methodology was successful, and we found that the same information can be obtained from the detector currents both for multiplicity counting [6], [7] and for reactivity measurements [8]. Experimental verification of the methodology is underway in measurements at the educational reactor of the Budapest University of Technology and Economics [9].

3. POWER REACTOR NOISE

3.1. Methodology

In power reactors several technological processes (boiling of the coolant/moderator in a BWR, vibrations of mechanical constructions in a PWR etc., commonly called “noise sources”), cause the medium in which the neutron transport and multiplication takes place to fluctuate in time and space. These fluctuations of the reactor material, expressed in terms of cross section fluctuations, will cause fluctuations of the neutron flux around its mean value. The measured neutron noise can be used either to detect, identify and quantify the noise sources, if the dynamic transfer function of the unperturbed core is completely known, or to monitor changes in the state

of the unperturbed core, such as reactivity coefficients, stability properties, if the properties of the noise source can be assumed to be known.

The methodology used to describe these stochastic processes is based on the linearised form of the Langevin technique. Starting with the space-time dependent diffusion equations in the form

$$\frac{1}{v} \frac{\partial \phi(\mathbf{r}, t)}{\partial t} = D \nabla^2 \phi(\mathbf{r}, t) + [\nu \Sigma_f(\mathbf{r}, t)(1 - \beta) - \Sigma_a(\mathbf{r}, t)] \phi(\mathbf{r}, t) + \lambda C(\mathbf{r}, t) \quad (5)$$

and

$$\frac{\partial C(\mathbf{r}, t)}{\partial t} = \beta \nu \Sigma_f(\mathbf{r}, t) \phi(\mathbf{r}, t) - \lambda C(\mathbf{r}, t), \quad (6)$$

splitting up the neutron flux and the cross sections into mean values and fluctuations, neglecting second order terms, subtracting the static equations and eliminating the delayed neutron precursors with a temporal Fourier transform, one arrives finally at the following generic equation for the neutron noise:

$$\tilde{\mathbf{L}}(\mathbf{r}, \omega) \delta \phi(\mathbf{r}, \omega) = S(\mathbf{r}, \omega). \quad (7)$$

Here the noise source $S(\mathbf{r}, \omega)$ is given, in terms of the fluctuations of the absorption cross sections, as $S(\mathbf{r}, \omega) = \phi_0(r) \delta \Sigma_a(\mathbf{r}, \omega)$, and the operator $\tilde{\mathbf{L}}(\mathbf{r}, \omega)$ corresponds to the unperturbed state of the system. With the Greens' function of equation (7), the neutron noise can be expressed as

$$\delta \phi(\mathbf{r}, \omega) = \int G(\mathbf{r}, \mathbf{r}', \omega) S(\mathbf{r}', \omega) d\mathbf{r}'. \quad (8)$$

The purpose of power reactor noise studies is that, by measuring the neutron noise, to identify either the noise source by knowing the the transfer function from calculations, and hence identify anomalies in the core, by inverting (8), or to determine some parameters of the transfer function, not affected by the noise source, for known forms of the noise source (such as a white noise background noise).

Examples of the first type (identifying anomalies/disturbances) include [2]

- determining the location (radial position) of vibrating control rods in PWRs
- detecting and quantifying core barrel vibration properties in PWRs
- measuring velocity of coolant flow in PWRs
- determining the position of a local thermal hydraulic instability in the core of a BWR
- detecting and classifying detector tube vibrations and impacting in BWRs
- determining two phase flow parameters in BWRs.

Examples of determining system parameters (properties of the transfer function) include

- determining of the moderator temperature coefficient (MTC) of PWRs;
- determining the margins of instability of a BWR.

One case of power reactor noise diagnostics will be mentioned here for illustration (several examples are given in [2]). Fig. 4 shows the location of a thermal hydraulic instability together with the detector positions used in the noise analysis-based localization of the instability (being due to an improperly seated fuel assembly), together with the position pointed out by the noise analysis based localization technique.

3.2 New developments

The current line of research is focused partly on the development of efficient computational methods for the fast and accurate calculation of the transfer function $G(\mathbf{r}, \mathbf{r}', \omega)$

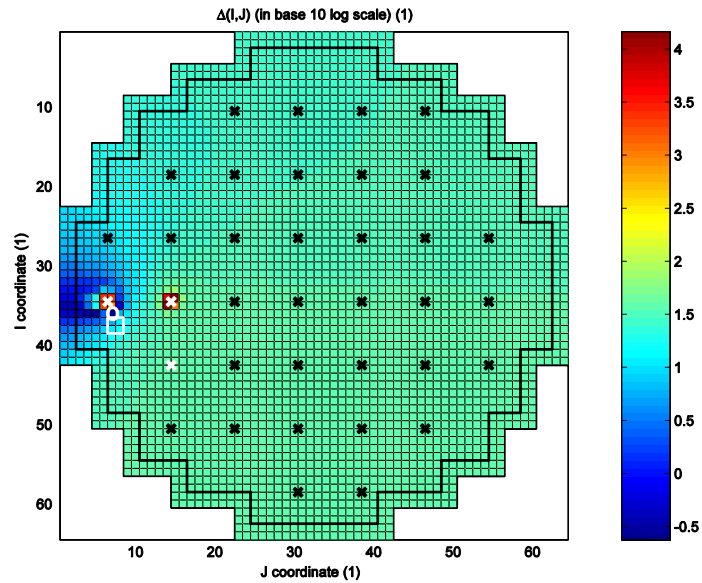


Fig. 2. Result of the localisation algorithm in the Forsmark-1 case (local instability event). The unseated fuel element is marked with a square, and the noise source identified by the localisation algorithm with a circle.

of real inhomogeneous cores, such as the dynamic core simulator CORE SIM [10], and inversion methods for unfolding the perturbation from the measured neutron noise by inverting Eq. (8) by machine learning methods. Another line of developments is the design of zero power reactor experiments, in support of the new models and diagnostic algorithms for power reactors. These lines are all pursued in the EU-supported project CORTEX [4]. Another current method under development is a methodology to determine the axial dependence of the void content along the whole height of a BWR from the measured neutron noise from the fix in-core detector signals. A neutron noise based method was currently suggested to solve this problem, and was tested in simulations [11]. Dedicated measurements are being set up in a special channel for inducing two-phase flow in the reflector of the zero power reactor CROCUS of EPFL with air injection (Fig. 3), to test the methodology and its performance, for a proof of principle. The experimental channel, and the two fission chambers next to it for cross-correlation measurements, are seen right to the core on the vertical cross section of the core (left hand side figure). This latter project is pursued in the project VOID of EPFL and PSI, funded by swissnuclear [12], [13].

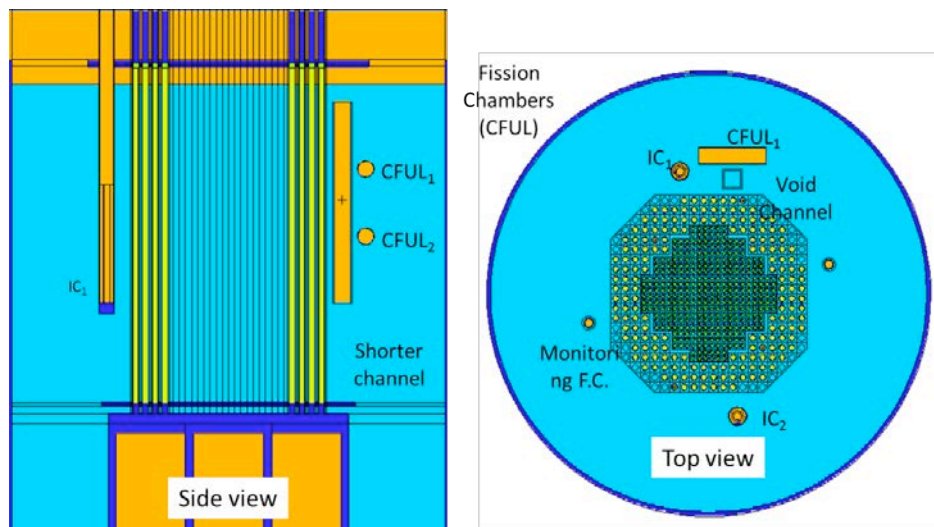


Fig. 3. A schematic view of the core of CROCUS and the experimental channel for air injection, with two fission chambers nearby for the measurement of the neutron noise.

4. CONCLUSION

The use of neutron fluctuations and neutron noise diagnostics is an active and intensively developing area of nuclear engineering and nuclear safeguards. In parallel with theory and method development, dedicated experiments and measurements performed in research reactors are a very effective tool to support research and development in both nuclear safeguards and power reactor diagnostics.

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